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SOLUTION OF CONVERGENT SEQUENCE BY USING THE HEAT DISTRIBUTION IN A THIN ROD

Heider Jabbar
College of Education

Udie Subrie
College of basic Education

Abstract

In this paper a studying for the heat equation in three dimensions is presented under boundary conditions and initial condition and through these conditions above, the solution for a convergence, in three dimension is performed.

Introduction

Many researchers studied the heat equation like [2],[3] and [4] because of the wide application of it in different fields. In this paper a study of heat equation in three dimension is considered with respect to the operators which effect on the accuracy of the solutions. We emphasize on the operators through the studying of the regular equations, corner equation and the boundary equations for a thin rod.

1. In this research we study the following problem

$$\varepsilon^2 \frac{\partial u}{\partial t} - a(\varepsilon^2 \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}) = \varepsilon^2 \sum_{|k| \leq m} f_k(u, x, t, \varepsilon) e^{ikwt}$$

$$(x, y, t) \in \Omega = (0 \leq x \leq 1) \times (0 \leq y \leq 1) \times (0 \leq t \leq T) \quad \dots (1)$$

$$u|_{t=0} = \varphi(x, y, \varepsilon), \quad u|_{x=0} = \sum_{|k| \leq m} \bar{\Psi}_{1k}(y, t, \varepsilon) e^{ikwt}$$

$$u|_{x=1} = \sum_{|k| \leq m} \bar{\Psi}_{2k}(y, t, \varepsilon) e^{ikwt}, \quad \frac{\partial u}{\partial y}|_{y=0} = 0, \quad \frac{\partial u}{\partial y}|_{y=1} = 0 \quad \dots (2)$$

ε is an arbitrary vector, $a = \text{const.} > 0$, $w, \varepsilon > 0$ a very small parameter. Later we shall discuss the problem in the region Ω , and we shall prove that the analytic convergent series

$$u(x, y, t, \varepsilon) = \sum_{i=1}^{\infty} \varepsilon^i u_i(x, y, t)$$

is a solution for the problem (1) with boundary conditions and initial condition in (2). In order to prove the existence of a solution for (1) and (2) we find an estimation for the convergent series, in other meaning for construct like this convergence and more accurate it contains two functions.

The first is the function the boundary of Ω when $t=0$, $x=0$, $x=1$, and the corner function at the points $(x,t)=(0,0)$ and $(x,t)=(1,0)$ in the region Ω .

2- The convergent series which is analytic with respect to, ε to solve (1), and (2) it must given as follows:-

$$u(x, y, t, \varepsilon) = \sum_{k=0}^{\infty} \varepsilon^k \left[u_k(x, y, t) + v_k(x, y, t, \tau) + B_k(x, y, \tau) + p_k(\zeta, y, t) + q_k(\zeta, y, t, \tau) \right. \\ \left. + r_k(\eta, y, t) + Z_k(\eta, y, t, \tau) + P_k(\zeta, y, \tau) + Q_k(\eta, y, \tau) \right] \quad \dots (3)$$

where u_k and v_k are the regular functions, B_k, P_k, q_k and r_k are boundary functions, and the corner functions P_k, Q_k such that

$\tau = \frac{t}{\varepsilon}, \xi = \frac{x}{\varepsilon}, \eta = (1-x)/\varepsilon$, moreover coefficient series (3) v_k, q_k and Z_k are with period 2π -periodic τ with zero average i.e.

$$\langle v_k(x, y, t, \tau) \rangle = \langle q_k(\xi, y, t, \tau) \rangle = \langle Z_k(\eta, y, t, \tau) \rangle = 0$$

$$\text{i.e. } \langle v_k(x, y, t, \tau) \rangle = \frac{1}{2\pi} \int_0^{2\pi} v_k(x, y, t, \tau) d\tau = 0$$

To find its coefficients (3) must be put in (1) and (2), we get the following equation

$$\frac{\partial u}{\partial \tau} - a(\varepsilon^2 \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}) = \varepsilon^2 \left(\sum_{i=0}^{\infty} \varepsilon^i \sum_{|k| \leq m} f_{ki}(u, x, t) e^{ik\tau} - \frac{\partial u}{\partial t} \right),$$

$$u|_{t, \tau=0} = \sum_{i=0}^{\infty} \varepsilon^i \phi_i(x, y), u|_{x=0} = \sum_{i=0}^{\infty} \varepsilon^i \sum_{|k| \leq m} \Psi_{1ki}(y, t) e^{ik\tau}$$

$$u|_{x=1} = \sum_{i=0}^{\infty} \varepsilon^i \sum_{|k| \leq m} \Psi_{2ki}(y, t) e^{ik\tau}, \frac{\partial u}{\partial y}|_{y=0,1} = 0 \quad \dots (4)$$

From the standard procedures for the analysis of series with respect to ε powers and by putting the coefficients for the equal powers with respect to ε , we get an equation for any convergent series

(a) The functions u_0 , and v_0 when ε^0 posses the following equations

$$a \frac{\partial^2 u_0}{\partial y^2} = 0, \frac{\partial u_0}{\partial y}|_{y=0,1} = 0 \quad \dots (5)$$

$$\frac{\partial v_0}{\partial \tau} - a \frac{\partial^2 v_0}{\partial y^2} = 0, \frac{\partial v_0}{\partial y}|_{y=0,1} = 0, \langle v_0(x, y, t, \tau) \rangle = 0 \quad \dots (6)$$

Which can be found $u_0 = \alpha_0(x, t)$, where α_0 -arbitrary function, and for the equation (6) we get $v_0 = 0$. By the same way we can find $u_1 = \alpha_1(x, t)$, where α_1 -arbitrary function, $v_1 = 0$.

For the coefficients u_2, v_2 , we get the following equations:

$$a \frac{\partial^2 u_2}{\partial y^2} = a \frac{\partial^2 \alpha_0}{\partial x^2} - \frac{\partial \alpha_0}{\partial t} + f_{00}(\alpha_0, x, t), \frac{\partial v_2}{\partial y}|_{y=0,1} = 0 \quad \dots (7)$$

$$\frac{\partial v_2}{\partial \tau} - a \frac{\partial^2 v_2}{\partial y^2} = \sum_{0 < |k| \leq m} f_{k0}(\alpha_0, x, t) e^{ik\tau} \quad \frac{\partial v_2}{\partial y}|_{y=0,1} = 0, \langle v_2(x, y, t, \tau) \rangle = 0$$

Decidability condition (7) we have the equation $\dots (8)$

$$\frac{\partial \alpha_0}{\partial t} - a \frac{\partial^2 \alpha_0}{\partial x^2} = f_{00}(\alpha_0, x, t) \quad \dots (9)$$

The solution of (8) evidently keep in mind

$$v_2 = \sum_{0 < |k| \leq m} \frac{1}{ik} f_{k0}(\alpha_0, x, t) e^{ikt}$$

Now coefficients $u_i, v_i, i \geq 3$ can be obtain problems:

$$a \frac{\partial^2 u_i}{\partial y^2} = a \frac{\partial^2 \alpha_{i-2}}{\partial x^2} - \frac{\partial \alpha_{i-2}}{\partial t} + \frac{\partial f_{00}(\alpha_0, x, t)}{\partial u} \alpha_{i-2} \quad \Phi_i(x, t) \frac{\partial u_i}{\partial y} \Big|_{y=0,1} = 0 \quad \dots(10)$$

where Φ_i can be expressed through the functions $v_s, s \leq i-3$ and $v_s, s \leq i-2$,

$$\frac{\partial^2 v_i}{\partial y^2} - a \left(\frac{\partial^2 v_{i-2}}{\partial x^2} + \frac{\partial^2 v_i}{\partial y^2} \right) = \sum_{0 < |k| \leq m} \Psi_{ki}(x, t) e^{ikt} \equiv \chi_i(x, t, \tau), \quad \frac{\partial v_i}{\partial \tau} \Big|_{y=0,1} = 0, \langle v_i(x, y, t, \cdot) \rangle = 0 \quad \dots(11)$$

where Ψ_{ki} can be expressed through the functions u_s and $v_s, s \leq i-2$. For example $\Phi_3(x, t) = f_{01}(\alpha_0, x, t)$,

$$\chi_3(x, t, \tau) = \sum_{0 < |k| \leq m} \left[\frac{\partial f_{k0}(\alpha_0, x, t)}{\partial u} \alpha_1(x, t) + f_{k1}(\alpha_0, x, t) \right] e^{ikt}$$

By decidability condition (10) we have the equation

$$\frac{\partial \alpha_{i-2}}{\partial t} - a \frac{\partial^2 \alpha_{i-2}}{\partial x^2} - \frac{\partial f_{00}(\alpha_0, x, t)}{\partial u} \alpha_{i-2} = \Phi_i(x, t) \quad \dots(12)$$

The solution of equation (11) becomes

$$v_i = \sum_{0 < |k| \leq m} \frac{1}{ik} \Psi_{ki}(x, t) e^{ikt}$$

(b) We find coefficients $P_k(\xi, y, t)$ and $q_k(\xi, y, t, \tau)$.

In order that consideration Eq. (1) in neighborhood left boundary at the region Ω , for that change of variables $\xi = x\varepsilon^{-1}$, lead to the problem

$$\frac{\partial u}{\partial \tau} - a \left(\frac{\partial^2 u}{\partial \xi^2} + \frac{\partial^2 u}{\partial y^2} \right) = \varepsilon^2 \left(\sum_{i=0}^{\infty} \varepsilon^i \sum_{|k| \leq m} f_{ki}(u, \varepsilon \xi, t) e^{ikt} - \frac{\partial u}{\partial t} \right) \quad u \Big|_{\xi=0} = \sum_{i=0}^{\infty} \varepsilon^i \sum_{|k| \leq m} \Psi_{1ki}(y, t) e^{ikt}, \frac{\partial u}{\partial y} \Big|_{y=0,1} = 0. \quad \dots(13)$$

By substituting series (3) in eq. (13) we get coefficients q_0 and p_0 as follows:

$$\frac{\partial^2 P_0}{\partial \xi^2} + \frac{\partial^2 P_0}{\partial y^2} = 0, \frac{\partial P_0}{\partial y} \Big|_{y=0,1} = 0, \quad (u_0 + P_0) \Big|_{x=0} = \Psi_{100}(y, t) \equiv \sum_{0 \leq s \leq m_00} a_s(t) \cos(s\pi y), P_0 \Big|_{\xi=\infty} = 0. \quad \dots(14)$$

$$\frac{\partial q_o}{\partial \tau} - a \left(\frac{\partial^2 q_o}{\partial \xi^2} + \frac{\partial^2 q_o}{\partial y^2} \right) = 0, \frac{\partial q_o}{\partial y} \Big|_{y=0,1} = 0,$$

$$q_o \Big|_{\xi=0} = \sum_{0 \leq |k| \leq m} \Psi_{lko}(y, t) e^{ikt} \equiv \sum_{0 \leq |k| \leq m} \sum_{0 \leq s \leq m_{ko}} b_{ks}(t) \cos(s\pi y) e^{ikt}, \quad \dots (15)$$

$$\langle q_o(\xi, y, t, \dots) \rangle = 0, q_o \Big|_{\xi=\infty} = 0$$

The solution of eq. (14) seek in the form:

$$P_o = \sum_{0 \leq |k| \leq m} C_s(\xi, t) \cos s\pi y,$$

in this connection we get unknown coefficients C_s from the following eq.:

$$\frac{d^2 C_s}{d\xi^2} - s^2 \pi^2 C_s = 0, C_o(0, t) = q_o(t) - \alpha_o(0, t),$$

$$C_s(0, t) = a_s(t) \text{ at } s \neq 0, C_s \Big|_{\xi=\infty} = 0$$

Then we find

$$P_o = \sum_{0 \leq |k| \leq m_{oo}} a_s(t) e^{-s\pi \xi} \cos s\pi y,$$

$$\alpha_o(0, t) = a_o(t) \equiv \langle \Psi_{100}(\cdot, t) \rangle.$$

The solution eq. (15) seek in the form:

$$q_o \Big|_{\xi=0} = \sum_{0 \leq |k| \leq m} \sum_{0 \leq |k| \leq m} C_{ks}(\xi, t) \cos \pi y e^{ikt}, \quad \dots (16)$$

Then

$$ik C_{ks} - a \left(\frac{\partial^2 C_{ks}}{\partial \xi^2} - \pi^2 S^2 C_{ks} \right) = 0, C_{ks}(0, t) = b_{ks}(t).$$

λ_{ks} can be expressed characteristic root of an equation

$a\lambda^2 = ik + a\pi^2 S^2$, at $\text{Re}\lambda > 0$. Thus we get

$$C_{ks} = b_{ks}(t) e^{-\lambda_{ks} \xi},$$

By virtue (16):

$$q_o \Big|_{\xi=0} = \sum_{0 \leq |k| \leq m} \sum_{0 \leq s \leq m_{ko}} b_{ks}(t) e^{-\lambda_{ks} \xi} \cos \pi y e^{ikt}$$

Similar, we find coefficient $P_i, q_i, i \leq 1$:

$$\frac{\partial P_{i-2}}{\partial t} - a \left(\frac{\partial^2 P_i}{\partial \xi^2} + \frac{\partial^2 P_i}{\partial y^2} \right) = \sum_{0 \leq s \leq m_i} M_{is}(t) P_{is}(\xi, \xi) e^{-M_{is} \xi} \cos \pi y \quad \dots (17)$$

$$\frac{\partial P_i}{\partial y} \Big|_{y=0,1} = 0, P_i \Big|_{\xi=0} = \Psi_{10i}(y, t) - \alpha_i(0, t), P_i \Big|_{\xi=\infty} = 0$$

$$\frac{\partial q_i}{\partial t} - a \left(\frac{\partial^2 q_i}{\partial \xi^2} + \frac{\partial^2 q_i}{\partial y^2} \right) = \sum_{0 \leq |k| \leq t_1} \sum_{0 \leq s \leq t_1} N_{iks}(t) Q_{iks}(\xi) e^{-N_{iks} \cos s \pi \tau} e^{ikt},$$

$$\frac{\partial q_i}{\partial y} \Big|_{y=0,1} = 0, \quad q_i \Big|_{\xi=0} = \sum_{0 \leq |k| \leq m_i} \Psi_{1kii}(y, t) e^{ikt} - v_i \Big|_{x=\infty} = 0 \quad \dots(18)$$

$$\langle q_i(x, y, t, \cdot) \rangle = 0, \quad q_i \Big|_{\xi=\infty} = 0$$

where

constant-

$$\sqrt{ni, ti, vi} > 0.$$

The solution eq. (17) seek in the form:

$$P_i \Big|_{\xi=0} = \sum_{0 \leq s \leq ni} C_{si}(\xi, t) \cos s \pi$$

The functions C_{si} are solution the following eq.

$$\frac{d^2 C_{is}}{d\xi^2} - s^2 \pi^2 C_{si} = M_{si}(t) P_{is}(\xi) e^{-M_{is} \xi},$$

$$C_{si}(0, t) = a_{si}(t) \text{ at } s \neq 0, \quad C_{si} \Big|_{\xi=\infty} = 0 \quad \dots(19)$$

which uniquely solution [3] and from equation $\alpha_i(0, t) = a_{oi}(t) - C_{oi}(0, t)$ we get formula

$$\alpha_i(0, t) \Big|_{t=0} = \langle \Psi_{10i}(\cdot, 0) \rangle. \quad \dots(20)$$

Now, we find the solution eq. (18). Let

$$\sum_{0 \leq |k| \leq m_i} \Psi_{1kii}(y, t) e^{ikt} - v_i \Big|_{x=\infty} = \sum_{0 \leq |k| \leq p_i} \sum_{0 \leq s \leq d_i} a_{ksj}(t) \cos s \pi y e^{ikt}$$

and

supposed

$\sqrt{t_i = p_i, v_i = d_i}$. The solution q_i will be seek in the form:

$$q_i = \sum_{0 \leq |k| \leq t_i} \sum_{0 \leq s \leq t_i} C_{ks}(\xi, t) \cos s \pi y e^{ikt}.$$

In this connection we get unknown coefficients C_{ks} form the following eq.:

$$ik - a \left(\frac{d^2 C_{ks}}{d\xi^2} - s^2 \pi^2 C_{ks} \right) = N_{iks}(t) Q_{iks}(\xi) e^{-v_{iks}},$$

$$C_{ks}(0, t) = a_{ksi}(t), \quad C_{ks} \Big|_{\xi=\infty} = 0$$

which it is obvious uniquely solution [2]. Coefficients r_k and Z_k are defined similar way. Moreover, in specifically we get:

$$\alpha_i(1, t) \Big|_{t=0} = \langle \Psi_{20i}(\cdot, 0) \rangle. \quad \dots(21)$$

(c) we find boundary function $B_k(x, y, \tau)$ in neighborhood lower boundary ($t=0$) at the region Ω and use change of variables $\tau = e^2 \tau$ in the problem (1), we get

$$\frac{\partial u}{\partial \tau} - a \frac{\partial^2 u}{\partial y^2} = \varepsilon^2 \left(a \frac{\partial^2 u}{\partial x^2} + \sum_{|k| \leq m} f_k(u, x, \varepsilon^2 \tau, \varepsilon) e^{ik\tau} \right) \quad \dots (22)$$

$$\frac{\partial u}{\partial y} \Big|_{y=0,1} = 0, u \Big|_{\tau=0} = \mathcal{G}(x, y, \varepsilon) = \sum_{|k| \leq m} \varepsilon^i \mathcal{G}_i(x, y).$$

From (3) and (22) function $b_0(x, y, \tau)$ have eq.

$$B_0(x, y, \tau) = \varphi(x, y, \tau) - u_0(x, y, \tau) - v_0(x, y, \tau, \tau). \quad \dots (23)$$

Taking into account $u_0(x, y, \tau) = \alpha_0(x, \tau)$ and

$v_0(x, y, \tau, \tau) = 0$ the eq. (23)

can be write

$$B_0(x, y, \tau) = \mathcal{G}_0(x, y) - \alpha_0(x, \tau) = \sum_{0 \leq s \leq n_0} b_{s0}(x) \cos s\pi y.$$

where n_0 -number and b_s -functions. The solution (23) seek in the form:

$$B_0(x, y, \tau) = \sum_{0 \leq s \leq n_0} C_s(x, \tau) \cos s\pi y.$$

and be found

$$C_0(x, \tau) = 0, C_s(x, \tau) = b_s(x) e^{-as^2 \pi^2 \tau}$$

$$B_0(x, y, \tau) = 0, B_0 \Big|_{\tau=\infty} = 0.$$

That mean $\alpha_0(x, \tau) = \langle \varphi_0(x, \tau) \rangle. \quad \dots (24)$

Similar, we have coefficient $B_1(x, y, \tau)$, from that we can be define

$\alpha_1(x, \tau) = \langle \varphi_1(x, \tau) \rangle$. Coefficients $B_i, i \geq 2$ series (3) are the following solutions equations:

$$\frac{\partial B_i}{\partial \tau} - \left(\frac{\partial^2 B_{i-2}}{\partial x^2} + \frac{\partial^2 B_i}{\partial y^2} \right) \sum_{0 \leq k \leq m_i} d_{ksi}(x) P_{ksi}(\tau) e^{ik\tau} e^{-\gamma_{ksi}\tau} \cos s\pi y \quad \dots (25)$$

$$B_i(x, y, \tau) = \mathcal{G}_i(x, y) - \alpha_i(x, \tau) - v_i(x, y, \tau, \tau), B_i \Big|_{\tau=\infty} = 0$$

where numbers $n_i, m_i, \gamma_{ksi} > 0$, functions d_{ksi} and P_{ksi} can be expressed through the functions B_s, v_s and $v_s, s \leq i-2, v_i(x, y_0, \tau) = 0$.

Let

$$\mathcal{G}_i(x, y) = \sum_{0 \leq s \leq n_i} d_{si}(x) P_{ksi} \cos s\pi y$$

Then we shall seek the solution equations (25) in the form:

$$B_i(x, y, \tau) = \sum_{0 \leq s \leq n_i} C_{si}(x, \tau) \cos s\pi y \quad \dots (26)$$

Thus we get for coefficients C_{si} the following equation:

$$\frac{\partial C_{si}}{\partial \tau} + as^2 \pi^2 C_{si} = \sum_{0 \leq l \leq n_i} d_{ksi}(x) P_{ksi}(\tau) e^{ik\tau} e^{-\gamma_{ksi}\tau}$$

$$C_{si} \Big|_{\tau=\infty} = 0, C_{si}(x, \tau) = C_{si}(x), 0 \leq s \leq m_i.$$

From that, we obtain uniquely functions $C_{si}(x, \tau)$ and taking into account (26), we define that

$$C_{oi}(X, 0)|_{x=0,1} = 0, \alpha_i(X, 0) = a_{oi}(X) - C_{oi}(X, 0),$$

$$\alpha_i(X, 0)|_{x=0,1} = \langle \vartheta_i(X, \cdot) \rangle_{x=0,1} \quad \dots(27)$$

(d) We observe that, by virtue equation (20), (21) and (27) are make consistent up to the continuity.

(e) We find coefficient P_k and Q_k . For P_o , as follow by using the mathematic induction.

$$\frac{\partial P_o}{\partial \tau} - a \left(\frac{\partial^2 P_o}{\partial \xi^2} + \frac{\partial^2 P_o}{\partial \xi^2} \right) = 0, \text{ at } k=0 \quad \dots(28)$$

$$P_o(\xi, \tau, 0) = -\rho_o(\xi, \tau, 0) - q_o(\xi, \tau, 0, 0), P_o(0, y, \tau) = B_o(0, y, \tau)$$

$$\frac{\partial P_o}{\partial y} \Big|_{y=0,1} = 0$$

According to (b) and (C) it should are representation

$$P_o(\xi, y, 0) = \sum_{0 \leq s \leq n_1} a_s(0) e^{-\lambda s \xi} \cos s \pi y,$$

$$B_o(0, y, \tau) = \sum_{0 \leq s \leq n_0} b_s(0) e^{-\lambda s \xi} \cos s \pi y$$

where $R_o \lambda s > 0$ and in this case, evidently, $P_o(0, y, 0) = B_o(0, y, 0)$. That means $m_{oo} = n_o, a_s(0) = b_s(0)$. The solution equation (28) we shall seek in the form:

$$P_o = \sum_{0 \leq s \leq o} N_s(\xi, \tau) - b_s(0) e^{-(\lambda s \xi + a s^2 \pi^2 \tau)} \cos s \pi y, \quad \dots(29)$$

Then for v_s we obtain the equation:

$$\frac{\partial v_s}{\partial \tau} - a \frac{\partial^2 v_s}{\partial \xi^2} + a \pi^2 s^2 v_s = -b_s(0) a \lambda_s^2 e^{-(\lambda s \xi + a s^2 \pi^2 \tau)} \equiv h_s(\xi, \tau), \quad \dots(30)$$

$$v_s(\xi, 0) = 0, v_s(0, \tau) = 0.$$

Introducing Greena function [2] we define

$$G_s = \frac{1}{2\sqrt{a\pi(\lambda - u)}} e^{-a\pi^2 s^2(\tau - \nu)} \left[e^{-\frac{(\xi - S)^2}{4a(\tau - u)}} - e^{-\frac{(\xi + S)^2}{4a(\tau - u)}} \right] \dots(31)$$

We produce solution equation (30) in the form [2]:

$$v(\xi, t) = \int_0^\infty du \int_0^\infty G_s(\xi, S, \tau - u) h_s(s, u) ds$$

for finding estimation function v we examine integral [3]:

$$J = \int_0^\infty du \int_0^\infty \frac{1}{2\sqrt{a(\tau - u)}} e^{-\gamma_1(\tau - u)} e^{-\frac{(\xi - S)^2}{4a(\tau - u)}} e^{-\gamma_2(u + s)} ds$$

where $\gamma_1, \gamma_2 > 0$. If we divide by substitution

$$\beta = \frac{\xi - S}{2\sqrt{a(\tau - \nu)}}, \text{ we get}$$

$$J = \int_0^\tau du \int_{\frac{\xi}{2\sqrt{a(\tau-\nu)}}}^{\infty} e^{-\gamma_1(\tau-u)} e^{-\beta^2} e^{-\gamma_2(u+\xi+2\sqrt{a(\tau-\nu)})} d\beta \leq$$

$$e^{-\gamma^2 \xi} \int_{-\infty}^{\infty} e^{-(\beta+\gamma_2\sqrt{a(\tau-\nu)})^2} d\beta \int_0^\tau e^{-(\gamma_1+\gamma_2 a)(\tau-\nu)} \cdot e^{-\gamma_2 \nu} d\nu \leq c e^{\gamma(\xi+\tau)}$$

where c=const., $\gamma = \min \frac{1}{2} (\gamma_2, \gamma_1 + \gamma_2^2 a)$. From relations (29), (31), (33) it should

be existence positives numbers C_0 and γ_0 have the following power estimation[4]:

$$|P_0(\xi, y, \tau)| \leq C_0 e^{-\gamma_0(\xi+\tau)}, \quad \xi, \tau \geq 0.$$

For functions $P_k, k \geq 1$, similar we find power estimation of boundary layer variables ξ and τ .

In the same way are construct equations for the corner boundary functions Q_k and have deduce power estimation.

References

- [1] Dieudonce J. Foundation of Modern analysis. Naok Moscow. 1984.
- [2] Levenshtam V.B. Asymptotic integration for differential equations with quick oscillatory composed Math. Cb. 2003. T. 81N:1357.
- [3] Levenshtam V.B., Abood H.D. Asymptotic integration of the problem on heat distribution in a thin rod with rapidly varying sources of heat. Contemporary Mathematics and it application. Vol. 7(2003). Moscow. 137-149.
- [4] Techanoff A.N. Samarecki A.A. Mathematical-physicist equation. Naok. Moscow. 1988.
- [5] Kreyszig E. Introductory functional analysis with applications. John Wiley and Sons. New York. 1990.
- [6] Smith D.R. Singular-Perturbation Theory. An Introduction with application. Cambridge Univ. Press Cambridge. 1985.