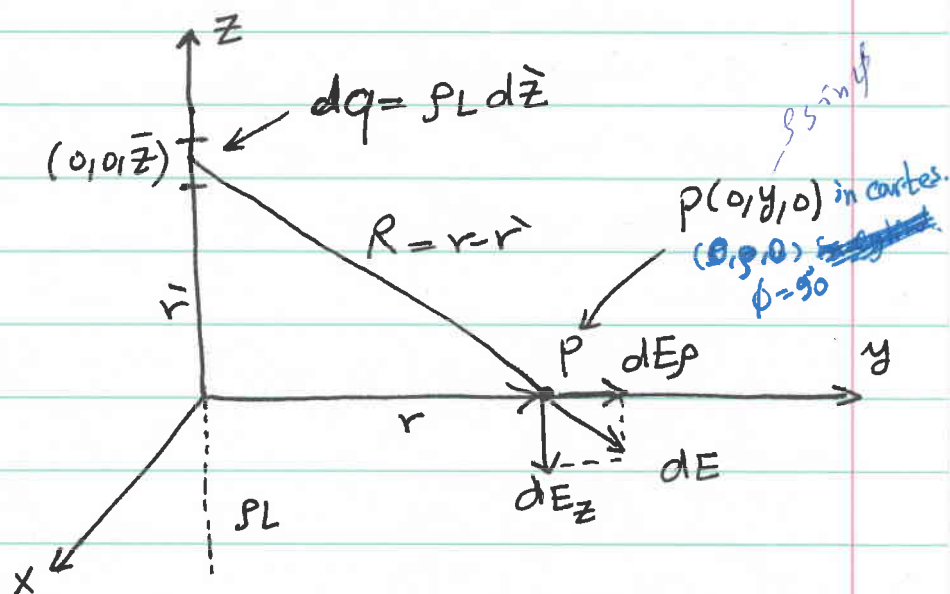


(11)

## Field of infinite line charge

Let us assume a straight line charge extending along the  $z$  axis in a cylindrical coordinate system from  $-\infty$  to  $\infty$ , We desire the electric field intensity  $E$  at any and every point resulting from a uniform line charge density  $\rho_L$ .



Symmetry should always be considered first in order to determine two specific factors

- (1) With which coordinate, the field does not vary and
- (2) Which component of the field are not present.

From testing Symmetry we conclude that:

- the field varies only with  $\rho$ .
- We have only an  $E_\rho$  component.
- the contribution of  $E_z$  will cancel.
- No element of charge produces  $E_\phi \Rightarrow E_\phi = 0$

(12)

$$dE = \frac{\rho_L d\vec{z} (r - r')}{4\pi\epsilon_0 |r - r'|^3}$$

where  $r = y a_y$   
 $\bar{r} = \bar{z} a_z$

and  $r - \bar{r} = \rho a_\rho - \bar{z} a_z$

Therefore,

$$dE = \frac{\rho_L d\vec{z} (\rho \bar{a}_\rho - \bar{z} a_z)}{4\pi\epsilon_0 (\rho^2 + \bar{z}^2)^{3/2}}$$

Since only the  $E_\rho$  component is present, we may simplify:

$$dE_\rho = \frac{\rho_L \rho d\bar{z}}{4\pi\epsilon_0 (\rho^2 + \bar{z}^2)^{3/2}}$$

and

$$E_\rho = \int_{-\infty}^{\infty} \frac{\rho_L \rho d\bar{z}}{4\pi\epsilon_0 (\rho^2 + \bar{z}^2)^{3/2}}$$

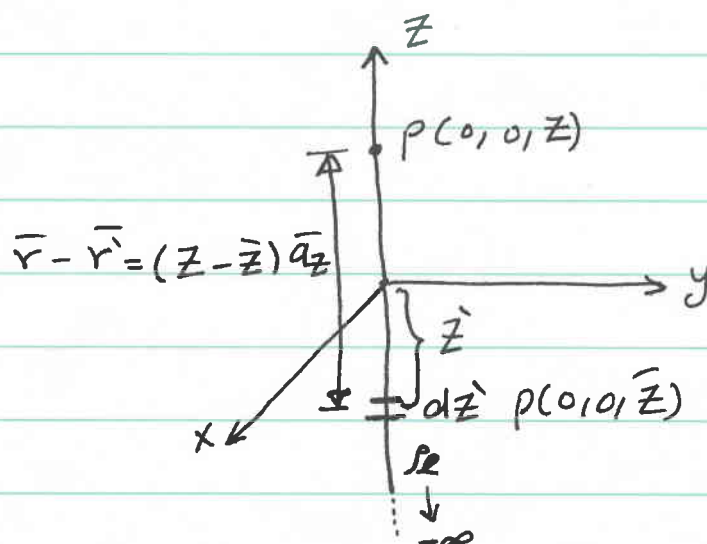
By integrating:

$$E_\rho = \frac{\rho_L}{4\pi\epsilon_0} \rho \left( \frac{1}{\rho^2} \times \frac{\bar{z}}{\sqrt{\rho^2 + \bar{z}^2}} \right)_{-\infty}^{\infty}$$

$$\therefore \boxed{E_\rho = \frac{\rho_L}{2\pi\epsilon_0 \rho}}$$

(B)

Example:- A semi-infinite line extending from  $-\infty$  to 0 along the  $z$ -axis carries a uniform charge distribution of  $100 \text{ nC/m}$ . Find the electric field intensity at point  $P(0,0,2)$ . If a charge of  $1 \text{ MC}$  is placed at  $P$ , calculate the force acting on it.



Solution:-

The differential charge element  $dq_i = \rho_l dz$  at  $z = \bar{z}$  from the origin.

The distance vector from  $\bar{z}$  to  $P$  is  $\vec{r} - \vec{r}' = (z - \bar{z}) \vec{a}_z$

The magnitude of the distance vector  $|\vec{r} - \vec{r}'| = z - \bar{z}$ .

$\therefore$   $E$  at point  $P$  is:

$$\vec{E} = \vec{a}_z \frac{\rho_l}{4\pi\epsilon_0} \int_{-\infty}^0 \frac{dz}{(z - z')^2}$$

$$\therefore \boxed{\vec{E} = \frac{\rho_l}{4\pi\epsilon_0 z} \vec{a}_z}$$

(124)

Substituting the values, we obtain

$$\vec{E} = \frac{9 \times 10^9 \times 100 \times 10^{-9}}{2} \vec{a}_z = 450 \vec{a}_z \text{ V/m}$$

The force acting on a charge of 1 MC at  $z = 2\text{m}$  is

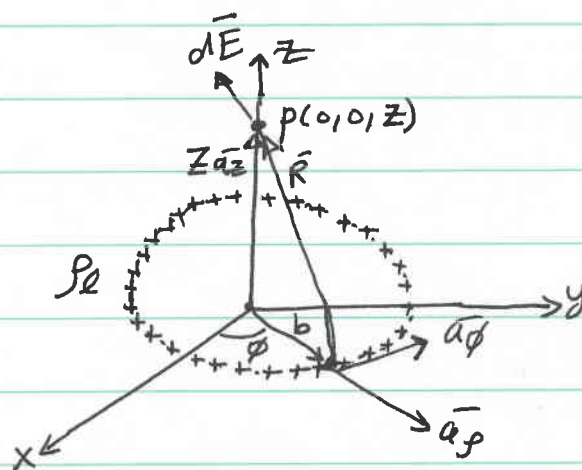
$$\vec{F} = q \vec{E} = 1 \times 10^6 \times 450 \vec{a}_z = 450 \vec{a}_z \text{ MN}$$

Example:- The charge is uniformly distribution in the shape of a ring of radius  $b$ , as shown in fig. Determine the electric field intensity at any point on the axis of the ring.

Solution:-

$$d\vec{q} = b d\phi \vec{a}_\phi$$

$$\vec{R} = -b \vec{a}_\rho + z \vec{a}_z$$



$$\vec{E} = \frac{\rho_L}{4\pi\epsilon_0} \int_0^{2\pi} \frac{b d\phi}{[b^2 + z^2]^{3/2}} (-b \vec{a}_\rho + z \vec{a}_z)$$

$$= \frac{\rho_L b}{4\pi\epsilon_0} \frac{1}{[b^2 + z^2]^{3/2}} \left[ -b \int_0^{2\pi} \vec{a}_\rho d\phi + z \int_0^{2\pi} d\phi \vec{a}_z \right]$$

Because  $\vec{a}_\rho = \vec{a}_x \cos \phi + \vec{a}_y \sin \phi$

$$\therefore \int_0^{2\pi} \vec{a}_\rho d\phi = \vec{a}_x \int_0^{2\pi} \cos \phi d\phi + \vec{a}_y \int_0^{2\pi} \sin \phi d\phi = 0$$

$$\therefore \vec{E} = \frac{\rho_L b z}{2\epsilon_0 [b^2 + z^2]^{3/2}} \vec{a}_z$$



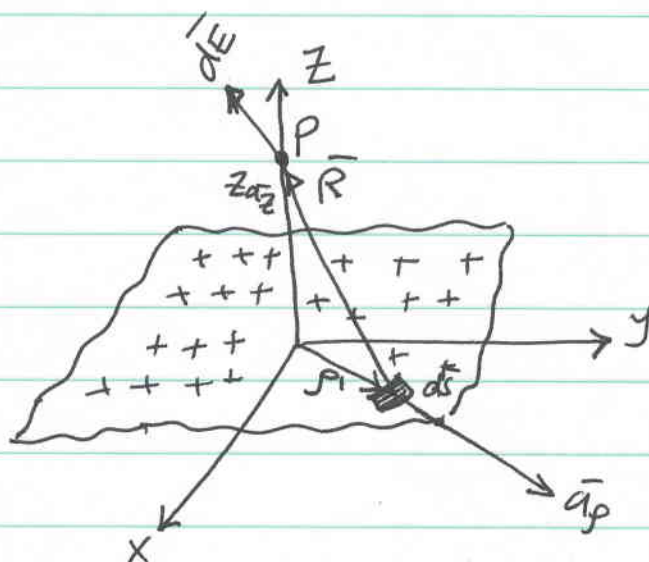
(15)

## Field of Infinite sheet of charges "

If charge is distributed with uniform density  $\rho_s$  (C/m<sup>2</sup>) over an infinite plane then to find  $\vec{E}$  at any point P the cylindrical coordinate system will be used with charge in the  $z=0$  plane.

Then

$$\vec{E} = \frac{\rho_s}{2\epsilon_0} \text{ an}$$



$$\vec{R} = -\rho \vec{a}_\rho + z \vec{a}_z \Rightarrow |\vec{R}| = \sqrt{\rho^2 + z^2}$$

$$dq = \rho_s ds \quad ds = \rho d\rho d\phi$$

$$\therefore dq = \rho_s \rho d\rho d\phi$$

$$\therefore d\vec{E} = \frac{\rho_s \rho d\rho d\phi}{4\pi\epsilon_0} \cdot \frac{-\rho \vec{a}_\rho + z \vec{a}_z}{(\rho^2 + z^2)^{3/2}}$$

Since, there is a symmetry about the  $z$ -axis, it results in cancellation of the radial ( $\rho$ ) component

(16)

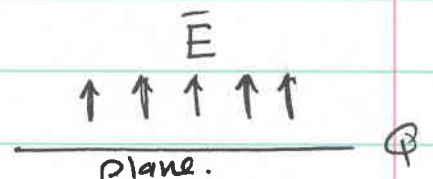
$$\vec{E} = \int_0^\infty \int_0^{2\pi} \frac{\rho_s \, r \, dr \, d\phi}{4\pi\epsilon_0 (r^2 + z^2)^{3/2}} z \vec{a}_z$$

$$= \frac{\rho_s}{4\pi\epsilon_0} \left[ \frac{-1}{\sqrt{r^2 + z^2}} \right]_0^\infty \left[ \phi \right]_0^{2\pi}$$

$$\Rightarrow \vec{E} = \frac{\rho_s}{2\epsilon_0} \vec{a}_z$$

This result is for points above xy-plane, below the xy-plane, the unit vector change to  $(-\vec{a}_z)$

The generalized form may be written using  $\vec{a}_n$  (the unit normal vector).

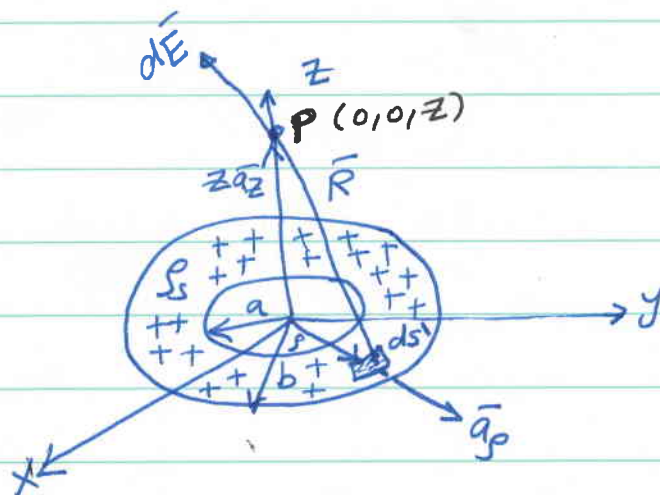
$$\boxed{\vec{E} = \frac{\rho_s}{2\epsilon_0} \vec{a}_n}$$


The electric field is everywhere normal to the plane of charge and its magnitude is independent of the distance from the plane.

(16)

(17)

Example: A thin annular disc of inner radius  $a$  and outer radius  $b$  carries a uniform surface charge density  $\rho_s$ . Determine the electric field intensity at any point on the  $z$ -axis when  $z > 0$ .



Solution:

$$dq = \rho_s \rho d\rho d\phi$$

$$\vec{R} = -\rho \vec{a}_\rho + z \vec{a}_z$$

$$|\vec{R}| = [\rho^2 + z^2]^{1/2}$$

∴ The electric field intensity at point  $P(0,0,z)$

is

$$\vec{E} = \frac{\rho_s}{4\pi\epsilon_0} \int_a^b \int_0^{2\pi} \frac{\rho d\rho d\phi}{[\rho^2 + z^2]^{3/2}} [-\rho \vec{a}_\rho + z \vec{a}_z]$$

$$\int_0^{2\pi} \vec{a}_\rho d\phi = 0$$

$$\vec{E} = \frac{\rho_s}{4\pi\epsilon_0} \int_a^b \int_0^{2\pi} \frac{\rho d\rho d\phi}{[\rho^2 + z^2]^{3/2}} z \vec{a}_z$$

$$= \frac{\rho_s z}{2\epsilon_0} \left[ \frac{1}{(a^2 + z^2)^{1/2}} - \frac{1}{(b^2 + z^2)^{1/2}} \right] \bar{a}_z \quad (8)$$

For an annular disc with very large outer radius  $b \rightarrow \infty$ , the electric field intensity becomes, fig 1

$$\bar{E} = \frac{\rho_s z}{2\epsilon_0} \left[ \frac{1}{(a^2 + z^2)^{1/2}} \right] \bar{a}_z$$

For a solid finite disc of outer radius  $b$ , the electric field intensity becomes, fig 2

$a=0$

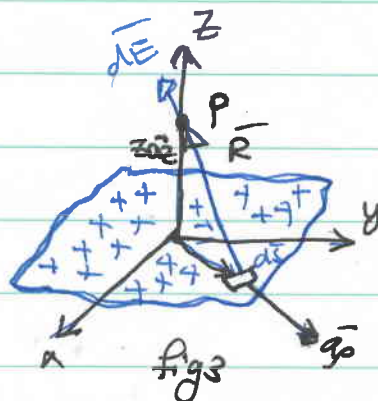
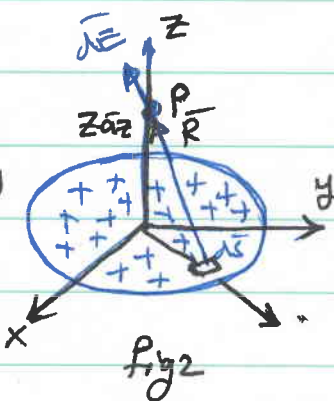
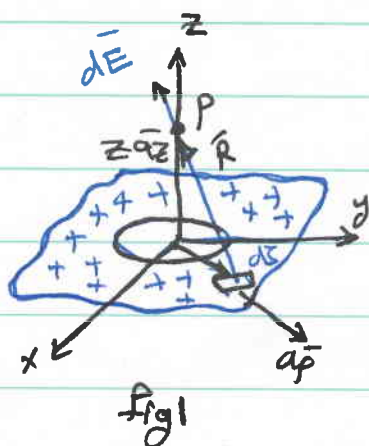
$$\bar{E} = \frac{\rho_s z}{2\epsilon_0} \left[ \frac{1}{z} - \frac{1}{(b^2 + z^2)^{1/2}} \right] \bar{a}_z$$

Finally the electric field intensity at any point due to an infinite plane of charge, fig 3

$a \rightarrow 0$   
 $b \rightarrow \infty$

Thus,

$$\bar{E} = \frac{\rho_s}{2\epsilon_0} \bar{a}_z$$





Example: Find the force on a point charge of  $50 \mu\text{C}$  at  $(0,0,5) \text{ m}$  due to a charge of  $500 \pi \mu\text{C}$  that is uniformly distributed over the circular disk  $r \leq 5 \text{ m}$ ,  $z=0 \text{ m}$

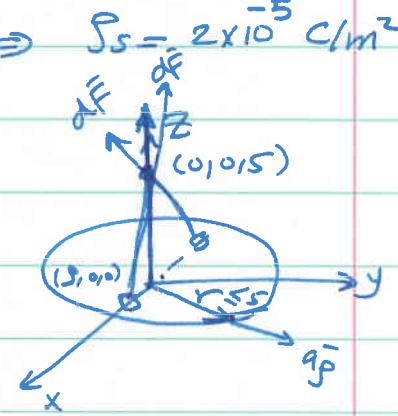
Solution:

$$\rho_s = \frac{Q}{A} = \frac{500 \pi \times 10^{-6}}{5^2 \times \pi} \Rightarrow \rho_s = 2 \times 10^{-5} \text{ C/m}^2$$

In cylindrical coordinates.

$$\vec{R} = -\rho \vec{a}_\rho + 5 \vec{a}_z$$

$$|\vec{R}| = \sqrt{\rho^2 + 25}$$



$$d\vec{F} = \frac{(50 \times 10^{-6}) (\rho_s d\rho d\phi)}{4\pi \epsilon_0} \cdot \frac{(-\rho \vec{a}_\rho + 5 \vec{a}_z)}{(\rho^2 + 25)^{3/2}}$$

Before Integration, Note that the radial components will cancel and that  $\vec{a}_z$  is constant.

$$d\vec{F} = 9 \times 10^9 \times 50 \times 10^{-6} (\overset{2 \times 10^{-5}}{\rho_s} d\rho d\phi) \cdot \frac{(5 \vec{a}_z)}{(\rho^2 + 25)^{3/2}}$$

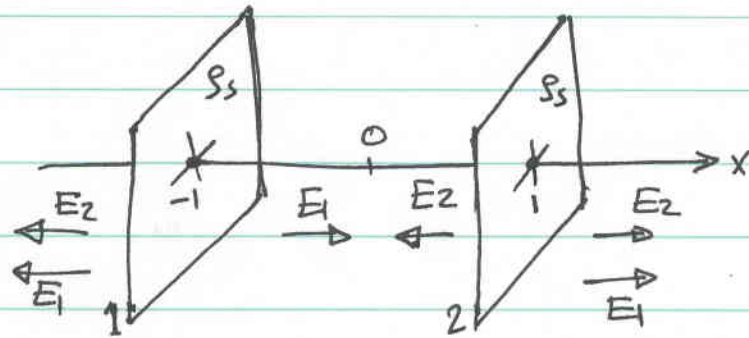
$$\therefore \vec{F} = 45 \int_0^5 \frac{\rho d\rho}{(\rho^2 + 25)^{3/2}} \int_0^{2\pi} d\phi$$

$$= 90\pi \left[ \frac{-1}{\sqrt{(\rho^2 + 25)}} \right]_0^5 \vec{a}_z$$

$$\therefore \vec{F} = 16.56 \vec{a}_z \text{ N.}$$

(20)

Example 8 Two infinite uniform sheets of charge, each with density  $\rho_s$ , are located at  $x = \pm 1$ . Determine  $\vec{E}$  in all regions.



Solution:

For region  $x > 1$

$$\vec{E}_1 = \frac{\rho_s}{2\epsilon_0} \vec{a}_x, \quad \vec{E}_2 = \frac{\rho_s}{2\epsilon_0} \vec{a}_x$$

$$\therefore \vec{E}_t = \vec{E}_1 + \vec{E}_2 = \frac{\rho_s}{\epsilon_0} \vec{a}_x \text{ (V/m)}$$

for region  $x < -1$

$$\vec{E}_1 = \frac{\rho_s}{2\epsilon_0} (-\vec{a}_x), \quad \vec{E}_2 = \frac{\rho_s}{2\epsilon_0} (-\vec{a}_x)$$

$$\therefore \vec{E}_t = \vec{E}_1 + \vec{E}_2 = \frac{\rho_s}{\epsilon_0} (-\vec{a}_x) \text{ (V/m)}$$

for region  $-1 < x < 1$

$$\vec{E}_t = \vec{E}_1 + \vec{E}_2 \Rightarrow \vec{E}_t = \frac{\rho_s}{2\epsilon_0} (-\vec{a}_x) + \frac{\rho_s}{2\epsilon_0} (\vec{a}_x) = 0$$

$$\therefore \vec{E}_t = \begin{cases} -\frac{\rho_s}{\epsilon_0} \vec{a}_x & x < -1 \\ 0 & -1 < x < 1 \\ \frac{\rho_s}{\epsilon_0} \vec{a}_x & x > 1 \end{cases}$$

(2)

HW. // Determine  $\vec{E}$  at  $(2, 0, 2)$  m due to three standard charge distributions as follows:-

- a uniform sheet at  $x=0$  m with  $\rho_{s1} = \left(\frac{1}{3\pi}\right) \text{ nC/m}^2$
- a uniform sheet at  $x=4$  m with  $\rho_{s2} = \left(-1/3\pi\right) \text{ nC/m}^2$  and
- uniform line at  $x=6$  m,  $y=0$  m with  $\rho_L = -2 \text{ nC/m}$ .

Ans:  $\vec{E} = 21 \vec{a}_x \text{ V/m}$ .

